Fuzzy Dark Matter

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- ✓ Background introduction
- □ Cosmology
- □ Galaxy
- □ Experimental studies

Outline of Background introduction

- What is Fuzzy Dark Matter (FDM)
- Particle physics motivation
- FDM numerical simulation
- Phenomenology

Before Fuzzy Dark Matter... What is Dark Matter?

- Existence of DM was first inferred by Zwicky in 1933, who discovered the "missing mass"
- Rich evidence for the existence of dark matter
- Range of possible mass is yet inconclusive

We are ignorant about the properties of DM

• Unknow range of possible mass & supporting theories



• Where should we consider DM as a wave instead of particle?

$$\lambda_{dB} = \frac{2\pi}{m\nu}$$

We are ignorant about the properties of DM

• Unknow range of possible mass & supporting theories



• Where should we consider DM as a wave instead of particle?

$$\lambda_{dB} = \frac{2\pi}{mv}$$
 gives m ~ 30 eV

FDM and CDM are similar in large scales

• FDM is similar to CDM in large scale





Schive, Chiueh, Broadhurst 2014

- In small scale FDM shows interference pattern, which motivates scientist on the cosmological formation of dwarf galaxy halos
- Get rid of the bothering issues of CDM simulation: smaller satellites problem, cusp-or-core problem, and "too big to fail" problem

Inconvenience of ultra-light particle

- Extremely hard to simulate
- CDM: adaptive mesh refinement (AMR) or smoothed particle hydrodynamics (SPH) techniques.
- FDM: you'll always need to resolve down to de Broglie wavelength in every cell of your box.



Particle physics motivations

Pseudo-Nambu-Goldstone boson



- Relic abundance sets the mass value
- Axion-like field with potential from non-perturbative effects

$$\mathcal{L} \sim -\frac{1}{2} (\partial \phi)^2 - \Lambda^4 (1 - \cos [\phi/F]) \qquad m \sim \Lambda^2/F$$

• Mis-alignment mechanism: mass range is relatively large



$$\Omega_{\text{matter}} \sim 0.1 \left(\frac{F}{10^{17} \,\text{GeV}} \right)^2 \left(\frac{m}{10^{-22} \,\text{eV}} \right)^{1/2}$$

Construction of wave halos

• Schrodinger-Poisson system

•
$$i\hbar \frac{\partial \psi}{\partial t} = \left(-\frac{\hbar^2}{2m_a}\nabla^2 + m_aV\right)\psi$$

•
$$\nabla^2 V = 4\pi G\rho = 4\pi G m_a |\psi|^2$$

Construction of wave halos -- fluid description

• Consider ψ as a classical fluid

•
$$\psi = \sqrt{\rho/m} e^{i\theta} \longrightarrow \rho = m|\psi|^2$$

• Mass conservation:
$$\frac{\partial \rho}{\partial t} + \nabla \cdot \rho v = 0$$
, where $v = \frac{1}{m} \nabla \theta$

• Euler equation:
$$\frac{\partial v}{\partial t} + v \cdot \nabla v = -\nabla V + \frac{1}{2m^2} \nabla \left(\frac{\nabla^2 \sqrt{\rho}}{\sqrt{\rho}} \right)$$

Numerical simulations

• Wave effects in a cosmological simulation.

$$\psi(t,\vec{x}) = \sum_{\vec{k}} A_{\vec{k}} e^{iB_{\vec{k}}} e^{i\vec{k}\cdot\vec{x} - i\omega_k t}$$



- Figure: a z = 5 snapshot of the dark matter density in a cosmological simulation of ultra-light dark matter with m = 10^{-22} eV
- Presence of interference fringes: a characteristic prediction of wave dark matter

Wave interference – vortices

- Naively, vorticity cannot exist because the velocity field is gradient flow.
- Consider fluid formation $\psi = \sqrt{\rho/m} e^{i\theta}$
- The loophole: when $\rho=0$
- No vortices in early universe
- Vortex generally takes the form of a loop i.e. vortex ring





Take-away messages

- FDM represents for the ultra-light DM, where we should consider as a wave instead of particles.
- FDM is more convenient in small-scale simulation compared to CDM, where you can see interference fringes pattern.
- The particle physics approach sets a range for FDM mass.
- There are 2 ways of wave construction of FDM: Schrodinger-Poisson system and fluid description.
- Existence of vortices is a important phenomenology of FDM.

Now let's talk about observational / experimental implications and constraints

Cosmology by Xiaochen Sun
Galaxy by Ruizhe Feng
Experimental studies by Jiejia Liu



Cosmology Implications & Constraints

Student Seminar



z~1100: CMB





$z \sim \mathcal{O}(1)$: (linear) matter power spectrum

• Fourier transfer from correlation function + Flux absorption from Ly α + Linear perturbation





$z \sim \mathcal{O}(1)$: (linear) matter power spectrum

- Fourier transfer from correlation function + Flux absorption from Ly α + Linear perturbation
- $P_m(k,z) = P_o(k) \times T^2(k,z)$
- How: quantum pressure
 - $p_a \propto -k^2 \rho \ln \rho / m_a$
 - pressure wins gravity on small scales

 \downarrow

suppress fluctuations





$z \sim \mathcal{O}(1)$: (linear) matter power spectrum

- Fourier transfer from correlation function + Flux absorption from Ly α + Linear perturbation
- $P_m(k,z) = P_o(k) \times T^2(k,z)$
- How: quantum pressure
- Uncertainty: IGM model & ionizing background
- Constraint: $m_a > 2 \times 10^{-20} \text{eV}$



Student Seminar



$z \sim \mathcal{O}(10)$: 21cm lines



- Uncertainty: stellar-to baryon fraction, minimum virial temperature & gas heating processes
- Constraint: $m_a > 8 \times 10^{-21} \text{eV}$ from EDGES



Mini Summary



• FDM always suppress small scale clustering due to quantum pressure.

Fuzzy dark matter: observational implications and constraints related to galaxies

Outline

- Explanation: Fornax dwarf galaxy
- Constraint: Subhalo mass function from stellar streams
- Prediction: Interference substructures

Fornax dwarf galaxy

- One of the Milky Way's neighboring dwarf galaxies
- It contains five globular clusters
- Dynamical friction should have caused
 - most of the clusters spiral to the center
 - merge to form a prominent nucleus
- BUT, this is NOT seen



Credit: ESO/Digitized Sky Survey 2

Dynamical friction

- Consider a subject body traveling through a population of field stars
- The field stars will be focused behind the subject body due to gravity, leading to a drag which slows down the subject body
- Dynamical friction will cause orbital decay of globular clusters



Credit: notes from Shude's course



Figure 8.3 The decay of the orbits of the Magellanic Clouds around our Galaxy.

From Binney & Tremaine, Galactic Dynamics

If the dark matter in Fornax is FDM rather than CDM ...

- The dynamical friction is suppressed
 - the wave nature of FDM is expected to suppress the overdensity, reducing the drag
 - standard estimates of the drag from dynamical friction must be modified
- Comparison between frictional decay times in FDM and CDM for the five Fornax clusters shows
 - Substantial increase in the time scale for dynamical friction
 - The orbital decay times are longer in a FDM halo than in a CDM halo
 - The shortest decay time in the FDM halo exceeds 2 Gyr

$$\tau = \frac{37.5 \text{ Gyr}}{C} \left(\frac{\mathfrak{M}(r)}{10^8 M_{\odot}} \frac{1 \text{ kpc}}{r} \right)^{3/2} \frac{10^5 M_{\odot}}{m_{\rm cl}} \frac{0.01 M_{\odot} \text{ pc}^{-3}}{\rho(r)}$$

	τ (Gyr)	τ (Gyr)
3×10 ⁻²² eV	112	215
	9.7	12
	0.62	2.2
	0.37	10
	21.3	31
	CDM	FDM

m =

Subhalo mass function

- The cold nature of DM causes it to form structure hierarchically starting from small gravitationally-bound halos that merge together to form larger halos
- The small halos are dense and highly concentrated, many of them survive tidal stripping and should exist as subhalos in galaxies today
- If the initial conditions of DM are warm (e.g. sterile neutrino models) or if DM is an ultra-light axion, then the abundance of lower-mass subhalos is significantly reduced

Stellar stream

- Elongated, almost one-dimensional structures produced by the tidal disruption of globular clusters or dwarf galaxies merging into the Milky Way
- They are highly sensitive to perturbations from passing DM subhalos
- Thus provide a means of measuring subhalo mass function



Stellar streams found in Gaia DR2



Observations of Pal 5 compared with simulations From: https://people.ast.cam.ac.uk/~derkal/files/pal5_pr/

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Power spectrum analysis

- ID power spectrum of the density contrast from the normalized linear density
 - which encodes the correlation in density contrast as a function of the angular scale
- Observations: GD-1 stream and Pal 5 stream
- Compare the observed power spectrum with the one predicted with mock realizations
 - Mock streams (baryonic substructures)
 - Mock streams + subhaloes (baryonic plus DM substructures from parameterized subhalo mass function)



FDM subhalo mass function is related to the particle mass

$$\left(\frac{dn}{dM}\right)_{\rm FDM} = \left[1 + \left(\frac{M}{2\,m_{22}^{-1.6} \times 10^8\,M_{\odot}}\right)^{-0.72}\right]^{-13.9} \left(\frac{dn}{dM}\right)_{\rm CDM} + \frac{0.014\,m_{22}^{1.5}}{M} \exp\left[-\left(\ln\left\{\frac{M}{4.7\,m_{22}^{-1.5} \times 10^8\,M_{\odot}}\right\}\right)^2/1.4\right] \right]$$



• Constraint: $m > 1.4 \times 10^{-21} \text{ eV}$





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Interference substructures

- A general feature of FDM models is the presence of ubiquitous density fluctuations
 - order unity density fluctuations on the scale of the de Broglie wavelength
 - can take the density all the way to zero (complete destructive interference, i.e., vortices)
 - distinct from subhalos as a form of halo substructure
 - these granular fluctuations occur because of the interference between bound waves in halos,
- These time-varying fluctuations can disturb the motions of stars, leading to potentially observable signatures in cold thin tidal streams in Galaxy
 - FDM model can generate significant small-scale structure in tidal streams

Calculation of FDM perturbations on tidal streams

- Mock stream: test particles placed uniformly along a circular ring at radius $r_0 = 15$ kpc
- The power spectra exhibit a sharp cutoff corresponding to the de Broglie wavelength of the FDM potential fluctuations
 - the gravitational forces are smooth on small scales, so the displacement perturbations and density perturbations are also smooth on small scales
- When stream perturbations become nonlinear, fold caustics generically arise
 - the cutoff disappears, and instead the power spectrum becomes a power-law
 - saturation in amplitude



(Dalal et al. 2021)

Take-home message

- FDM model can explain the survival of globular clusters against orbital decay in Fornax dwarf galaxy
- The subhalo mass function can be measured from perturbations of stellar streams, and put constraints on the particle mass of FDM
- The interference substructures predicted by FDM model can be probed by power spectrum of tidal streams, which will show a cutoff and fold caustics

Observational/Experimental Detection of FDM

Outlines

- Astrophysical: pulsar timing array
- Experimental:
 - CAPSEr-Wind
 - Atomic clock

Time dependent gravitational oscillations

Scalar field of FDM $\phi(\mathbf{x},t) = A(\mathbf{x}) \, \cos(mt + lpha(\mathbf{x}))$



Measure the time residual

Sensitivity

similar to a gravitational wave background with characteristic strain

$$h_c = 2\sqrt{3}\Psi_c = 2\cdot 10^{-15} igg(rac{
ho_{
m DM}}{0.3 {
m GeV/cm^3}}igg) igg(rac{10^{-23} {
m eV}}{m}igg)^2$$
 $f \equiv 2\pi\omega = 5\cdot 10^{-9}~{
m Hz}igg(rac{m}{10^{-23} {
m eV}}igg)$



Dimensionless Ψ_c

Sensitivity



Take home message

- Scalar field of FDM causes time delay of pulsar signal
- It is similar to signals caused by GW but has 2 major differences:
 - FDM signal is independent on the direction to pulsar
 - It is monochromatic and can appear as an excess in the signal at particular frequency $\propto m$
- Currently 🔀, future SKA 🔽

Axial nuclear moment

APL interaction with SM fermion:

 $\frac{\partial_{\mu}a}{f_a} \bar{\Psi}_f \gamma^{\mu} \gamma_5 \Psi_f$ APL nuclear current coupling



additional perturbation term in Hamiltonian

Energy oscillation amplitude

$$egin{array}{lll} H_N \supset g_{aNN}
abla a \cdot m{\sigma}_N \ \supset g_{aNN} m_a a_0 \cos(m_a t) {f v} \cdot m{\sigma}_N & \longrightarrow & \Delta E \sim g_{aNN} \sqrt{
ho_{DM} v} \end{array}$$

CAPSEr-Wind

CAPSEr (Cosmic Axion Spin Precession Experiment)

The lab has a relative motion with respect to the FDM wind

The FDM-nucleon interaction causes the spin to precession around *v*

The spin precession both around v and B

Transverse magnetization

$$M(t) \propto \sqrt{
ho_{DM}} \ rac{\sin[(2\mu B_{
m ext}-m_a)t]}{2\mu B_{
m ext}-m_a}$$



sensitivity issues

The detection in particle physics is model dependent, experiment sensitivity may reach the detection limit when FDM model is well-constructed



Atomic clock

$$\mathcal{L} \supset rac{1}{2} (\partial_\mu \phi)^2 - rac{1}{2} m_\phi^2 \phi^2 + rac{-1 + d_e \kappa \phi}{4 e^2} F_{\mu
u} F^{\mu
u} \quad riangleq lpha(t) \simeq lpha [1 + d_e \kappa \phi_0 \cos(m_\phi t + \delta)]^2$$



Transition frequency shift between A/B states:

- $v_{162} v_{164}$: α sensitive, measures the signal
- $v_{162} v_{164}$: α insensitive, correct systematic error



Atomic clock





Take home message

- Axion-nucleon interaction causes spin precession around the relative velocity
- transverse magnetization happens when applied magnetic field equals to the precession frequency, which $\propto m$
- To reach the limit of FDM is challenging using CAPSEr-Wind method (unless you change the model), atomic clock seems to have hope
- Some paper claims FDM can change the neutrino oscillation probability, and thus they can be detected in neutrino experiment.