

# Particle Acceleration in Astrophysics

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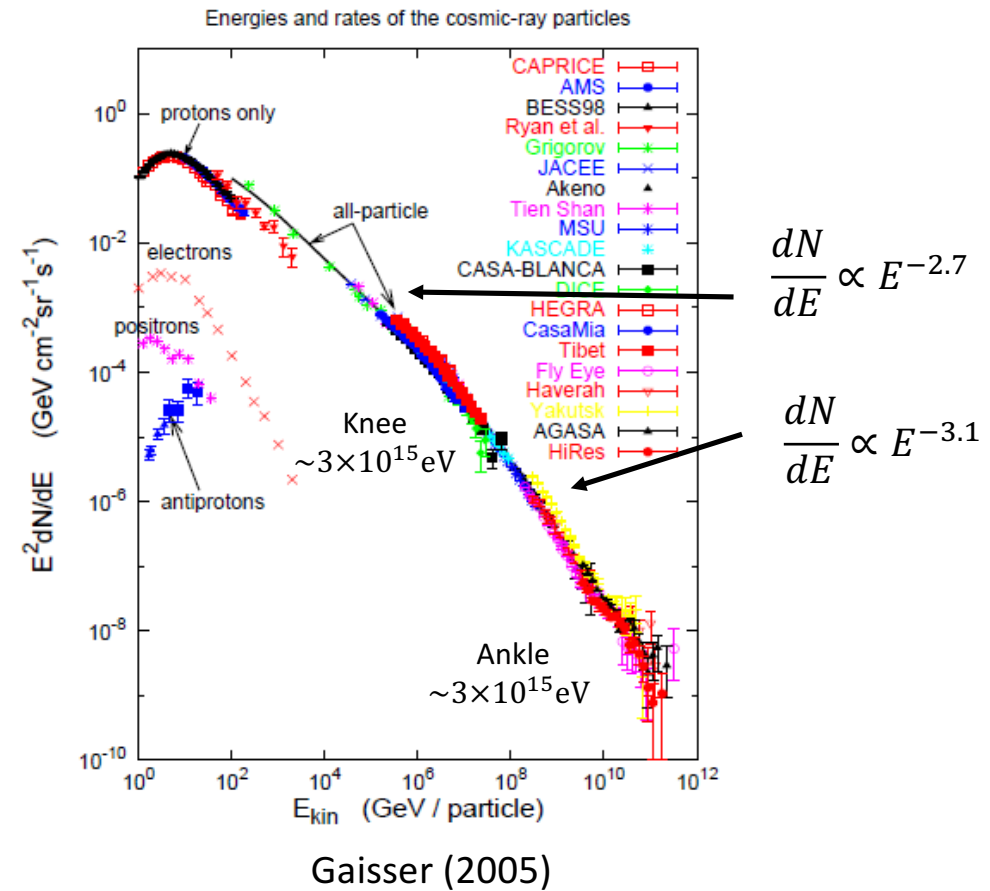
1. Second Order
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# Cosmic Rays (CRs): Introduction

- CRs refer to the **high-energy extraterrestrial particles**.
- CRs contribute an energy density of  **$1\text{eV}/\text{cm}^3$** .
- CRs are essentially **collisionless**.
- Most CRs drift and diffuse in **magnetic field**.



# CRs: Early Research

- Victor Hess (1912): Measured the ionizing radiation in the atmosphere thorough a series of balloon experiments and concluded the radiation came from **out of the earth**.
- Robert Millikan (1925): Named it “**Cosmic Rays**” .
- Jacob Clay (1927): Confirmed CRs were **charged particles**.
- **Where are they from?**
- Baade and Zwicky (1934): First proposed CRs originated from **supernova remnants (SNRs)**.



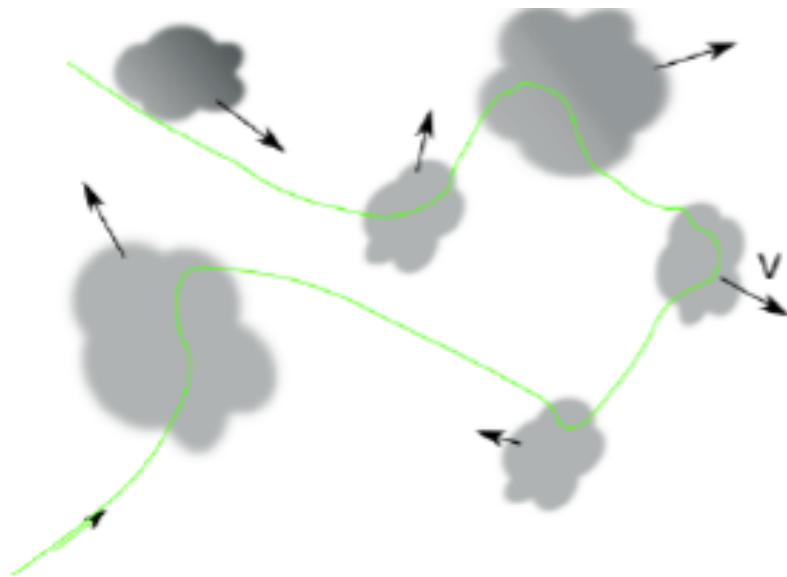
Hess (center) on the balloon (wikipedia)

# CRs Acceleration: General Considerations

- Lorentz Force:  $\frac{d(\gamma v)}{dt} = \frac{q}{m} (E + \frac{v}{c} \times B)$
- Magnetic Fields do not work!
- $E_{\perp} = -\frac{v_{\text{flow}}}{c} \times B$
- Only moving fields can work!
- $E_{\parallel}$
- Zero in ideal MHD!

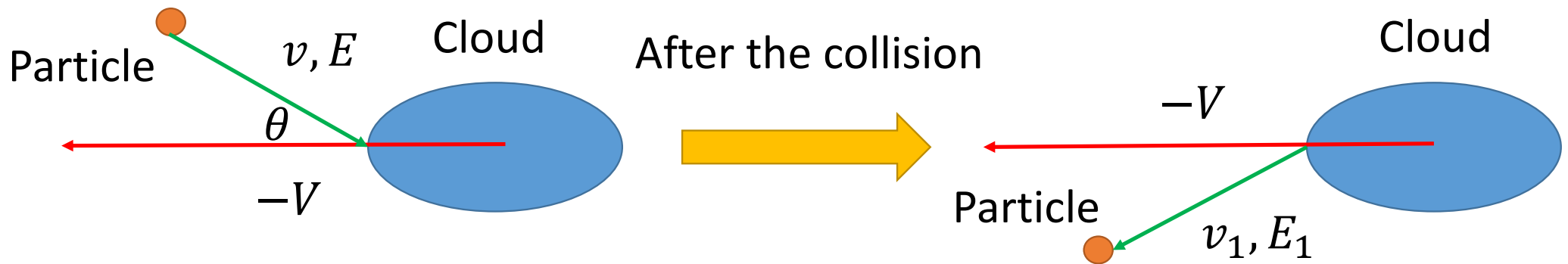
# Towards the Origin of CRs: Fermi Acceleration

- Enrico Fermi (1949): Particles are accelerated in ISM thorough **collisions** with “**moving magnetic fields**”.



An illustration for Fermi's model.  
(The magnetic fields are carried by **non-relativistic moving clouds** in ISM .  
CRs randomly scatter from the clouds and **statistically gain energy.** )

## Second-Order Fermi Acceleration :



- $E_1 = \gamma_V^2 E \left( 1 + \frac{2Vv \cos \theta}{c^2} + \frac{V^2}{c^2} \right);$
- $\frac{E_1 - E}{E} = \frac{\Delta E}{E} \approx \frac{2Vv \cos \theta}{c^2} + \frac{V^2}{c^2}$  (Expand to 2<sup>nd</sup> order of  $\frac{V}{c}$ );
- $f(\theta) d\theta = \frac{1}{2} \left( 1 + \frac{V \cos \theta}{c} \right) \sin \theta d\theta$  (The distribution function of  $\theta$ );
- $\left\langle \frac{\Delta E}{E} \right\rangle \approx \frac{8V^2}{3c^2}$

# Energy Spectrum:

- CRs gain energy at a rate that is **proportional to their energy**, and escape from the acceleration region in a Poisson process with **energy-independent probability**.
- Two scale time  $\tau_{acc}$  and  $\tau_{esc}$ .
- Initially,  $N_0$  particles with energy  $E_0$ ;
- After time  $t$ ,  $E = E_0 e^{t/\tau_{acc}}$ ,  $N = N_0 e^{-t/\tau_{esc}} = N_0 \left(\frac{E}{E_0}\right)^{-\frac{\tau_{acc}}{\tau_{esc}}}$ ;
- $\frac{dN}{dE} \propto E^{-\left(1+\frac{\tau_{acc}}{\tau_{esc}}\right)}$  **Power-law energy spectrum!**

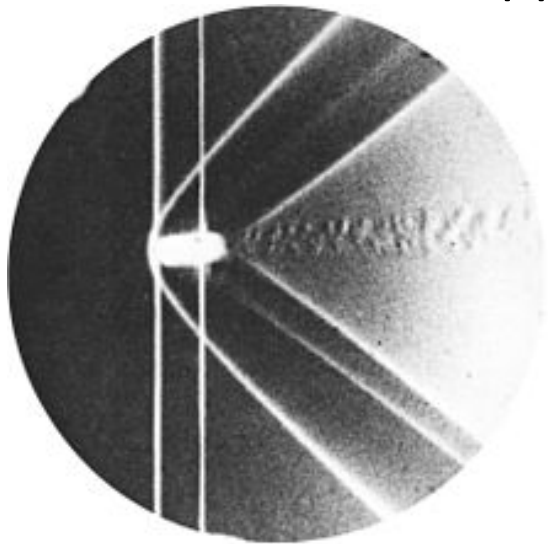


## Disadvantages of 2nd-order Fermi Acceleration:

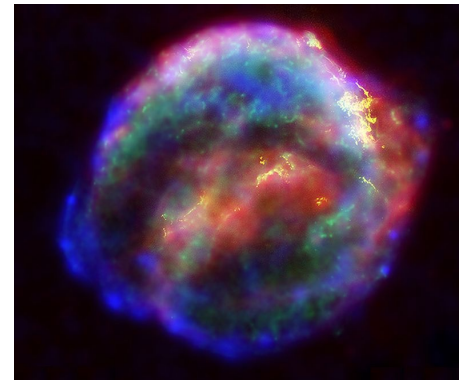
- Normally  $\frac{V}{c} \sim 10^{-4}$ , so 2nd-order Fermi Acceleration is **too slow** to drive particles to high energy.
- The process can produce a power-law energy spectrum, but the power-law index is **unconstrained**.
- There exists **injection threshold energy** for particle to overcome ionization loss.

# First-Order Fermi Acceleration: Diffusive Shock Acceleration (DSA)

- Bell(1978);Blandford & Ostriker (1978) proved the 1<sup>st</sup> order Fermi Acceleration could happen at the **shock front**.

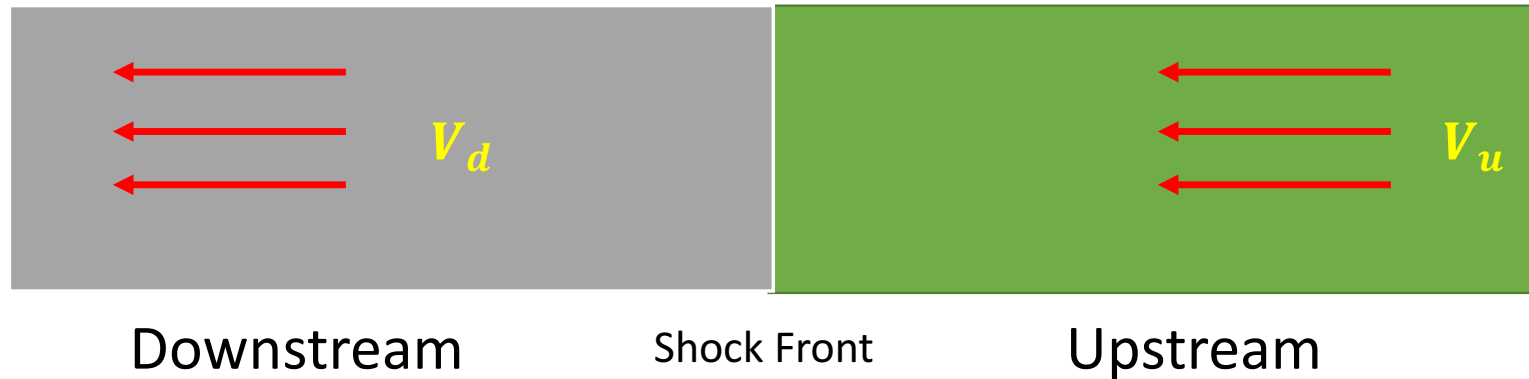


Shocks are formed when the perturbation in a fluid propagates faster than the sonic speed. Here is a picture of the shock on a flying bullet (E.Mach & P. Salcher 1887).



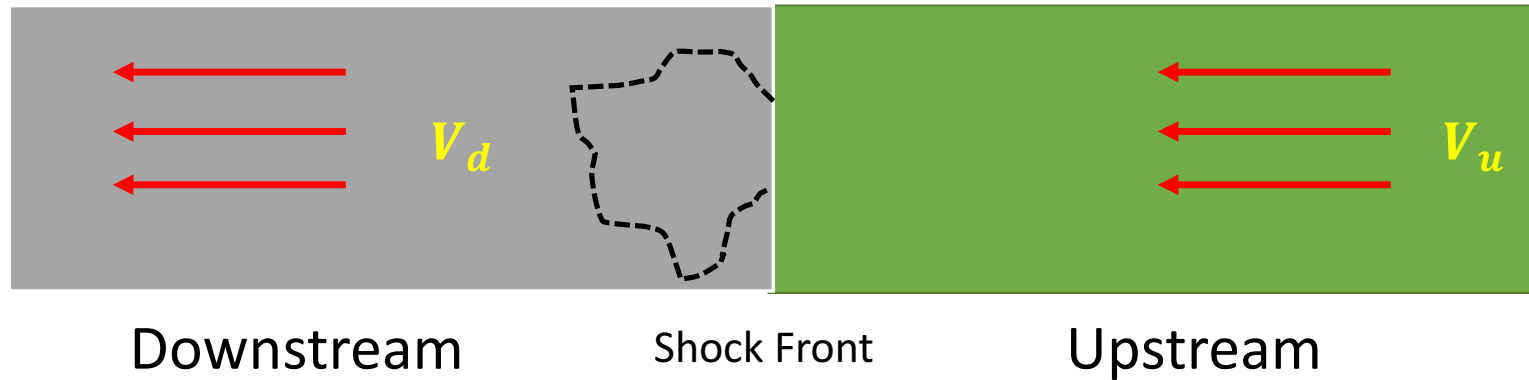
Shocks are common in astrophysics environment, e.g. supernova explosion and Gamma ray burst. Left: SNRs; Right:GRBs. (wikipedia)

# Physics of Shock:



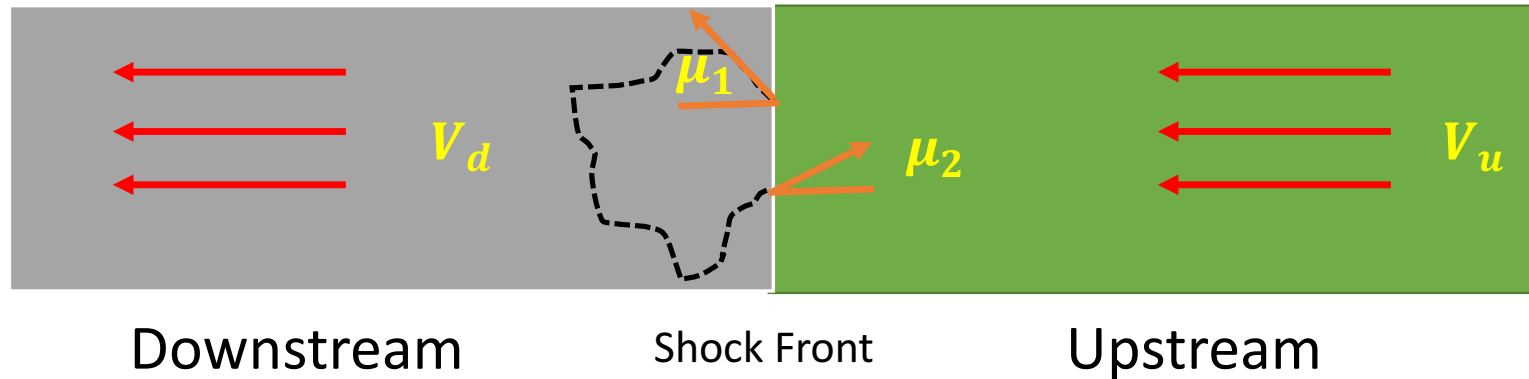
- Discontinuity:  
e.g. In high-Mach limit,  $\frac{\rho_d}{\rho_u} = \frac{v_u}{v_d} = r = 4$  ( $r$  is the compression ratio)
- Converging Flow
- Energy Dissipation

# Physics of Shock:



- Particles diffuse due to the turbulence.

# Physics of Shock:



- Energy initially in upstream frame :  $E_1$ ;
- Energy in downstream frame :  $E_2 = \gamma_V E_1 (1 + \beta \mu_1)$ ;
- Energy finally in upstream frame:  $E'_1 = \gamma_V E_2 (1 - \beta \mu_2)$ ;
- $\left\langle \frac{\Delta E}{E_1} \right\rangle = \frac{4}{3} \beta$
- $\beta = (V_d - V_u)/c, \gamma_V = (1 - \beta^2)^{-1/2}$

# Energy Spectrum:

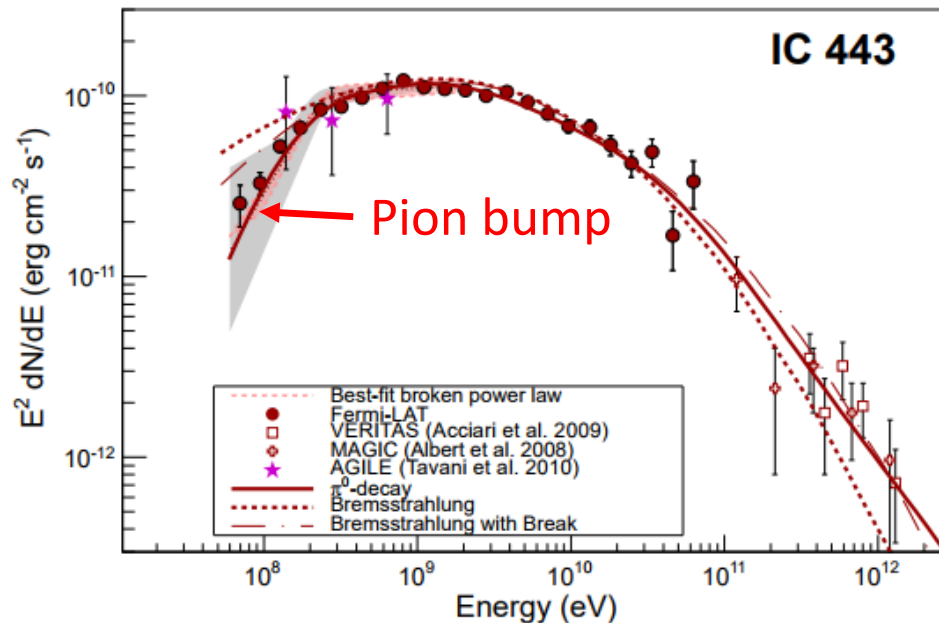
- Define the probability of CRs remain in accelerated region after one cycle:
- $J_+$ : the flux of CRs entering the shock from upstream
- $J_-$ : the flux of CRs returning to upstream from downstream
- $J_\infty$ : the flux of CRs escaping in far downstream
- $P = \frac{J_-}{J_+} = \frac{J_-}{J_- + J_\infty}$
- Index for power-law:  $s = 1 + \frac{3}{r-1}$  ( $s = 2$  when  $r = 4$ )
- From convection-diffusion equation:
  - $f(p) \propto p^{-4} \rightarrow f(E) \propto E^{-1.5}$  (non-relativistic particle)
  - $f(E) \propto E^{-2}$  (relativistic particle)

# Acceleration Rate:

- In one acceleration cycle for CRs (upstream-downstream-upstream):
- The duration time:  $t_{\text{cycle}} = \frac{4D_u}{V_u c} + \frac{4D_d}{V_d c}$ ;
- $D_u$  and  $D_d$  is the diffusion coefficient in upstream or downstream;
- Acceleration rate:  $\frac{dE}{dt} = \frac{4(V_d - V_u)}{3c} \frac{E}{t_{\text{cycle}}}$ ;
- In Bohm limit,  $\frac{dE}{dt} = \text{const} \rightarrow E_{\text{max}} \propto t$
- Without radiation loss, **the total energy increases linearly with time!**

# Maximum energy achieved in SNRs:

- For a SNR shock,  $t \sim 10^3 \text{ years}$ ,  $V_u \sim 5000 \text{ km s}^{-1}$ ,  $B_u \sim 10 \mu\text{G}$ ,
- The estimated maximum energy is  $\sim \text{PeV} (10^{15} \text{ eV})!$
- For higher energy, we should consider other sources!

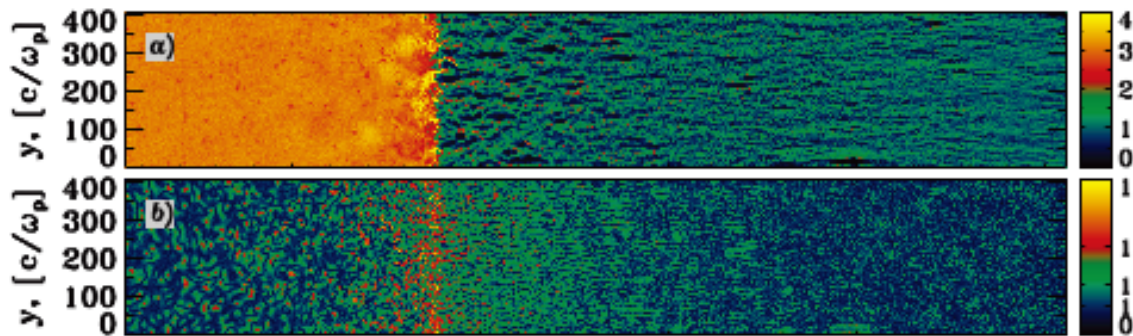


CRs acceleration in SNRs can be tested by e.g. the generated  $\gamma$ -ray emission. (Energy Spectrum of SNR IC443, M. Ackermann et al. 2013 )

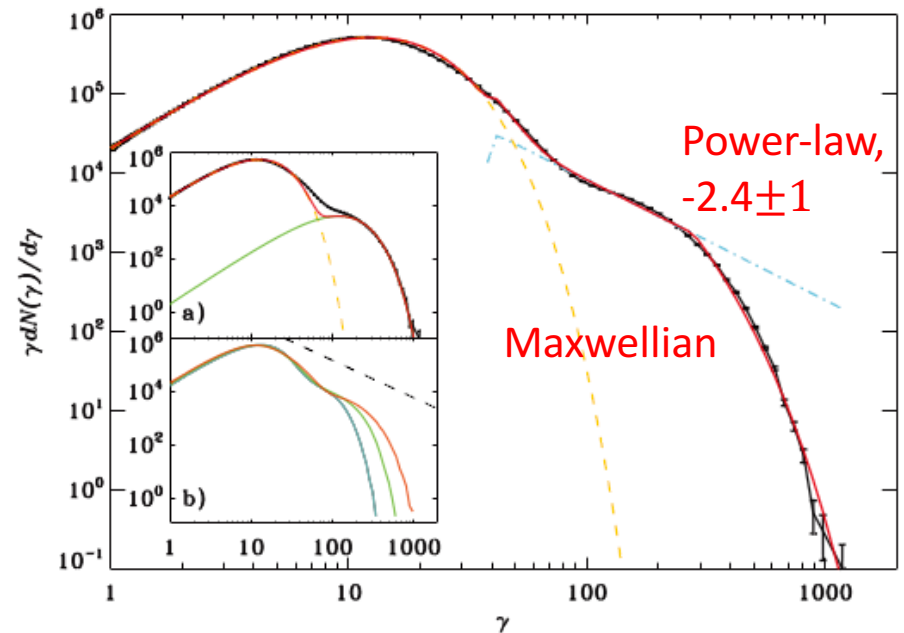


# Numerical Simulations (1):

A. Spitkovsky (2008) first simulated a **self-consistent** particle acceleration process with collisionless **relativistic** shock in **electron-positron** pair plasmas



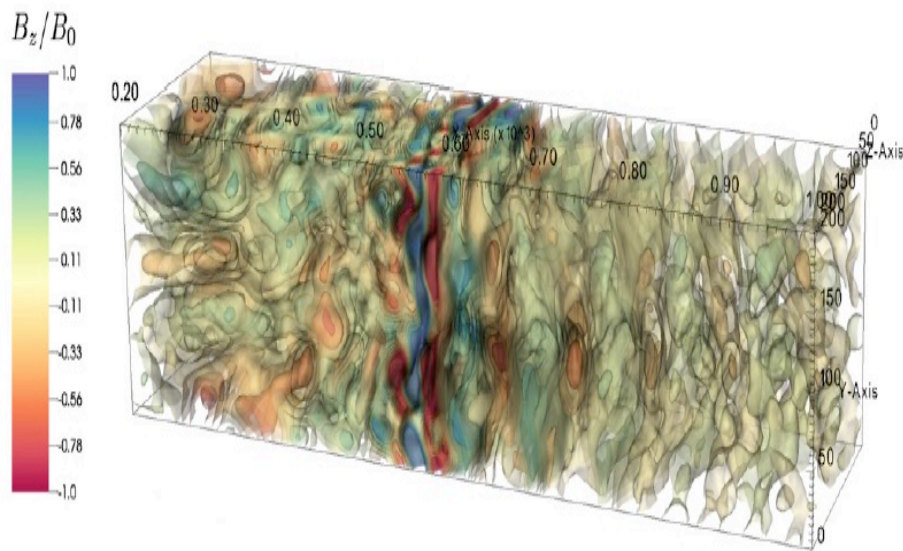
Density and magnetic energy density on the simulation plane (Spitkovsky 2008)



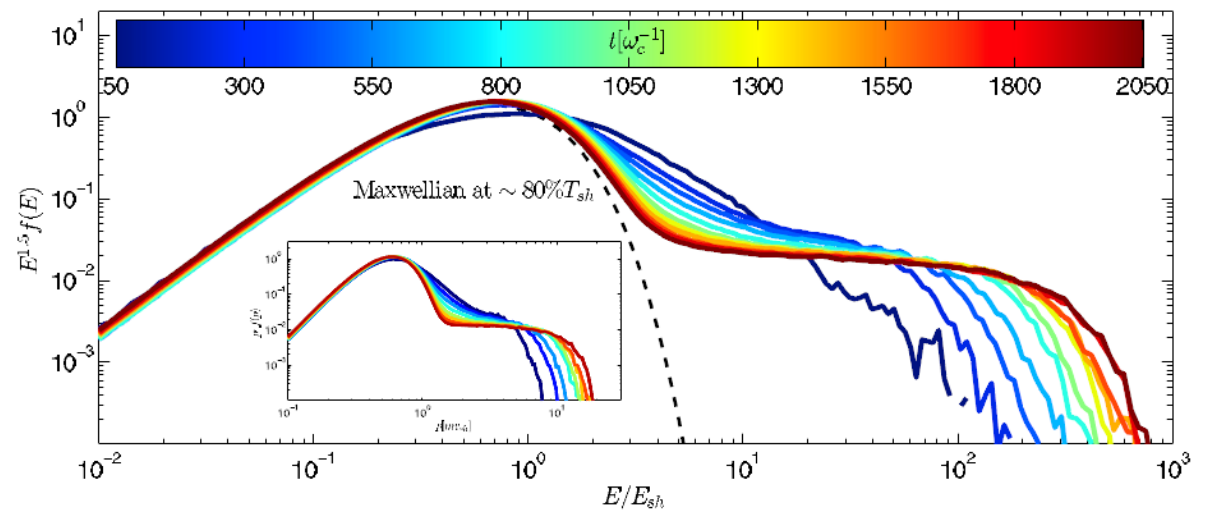
Particle energy spectrum (Spitkovsky 2008)

## Numerical Simulations (2):

- D.Caprioli & A. Spitkovsky (2013): simulated the self-consistent particle acceleration process with **non-relativistic** shock.



Self-generated Magnetic fields in the simulation box, (D.Caprioli & A. Spitkovsky 2013)



Downstream ion energy spectrum at different times (D.Caprioli & A. Spitkovsky 2013)

# Summary:

- Cosmic Rays are high energy extraterrestrial particles and important in astrophysics.
- Diffusion shock acceleration can explain the power-law properties of CRs energy spectrum.
- Most CRs are thought to be come from SNRs, and there have been some evidence.
- Present numeric simulations can have a self-consistent acceleration process.

# References:

1. E. Fermi(1949), On the Origin of the Cosmic Radiation.
2. R. Blandford & D. Eichler(1987), Particle Acceleration at Astrophysical Shocks: a Theory of Cosmic Ray.
3. A. Spitkovsky (2008), Particle Acceleration in Relativistic Collisionless Shocks: Fermi Process at Last?
4. M. Ackermann et al.(2013), Detection of the Characteristic Pion-decay Signature in Supernova Remnants.
5. P. Blasi (2013), The Origin of Galactic Cosmic Rays.
6. D. Caprioli & A. Spitkovsky (2013), Simulations of Ion Accelerations at Non-relativistic Shocks I. Acceleration Efficiency.
7. R. A. Treumann (2013), Fundamentals of Collisionless Shocks for Astrophysical Application, 1. Non-relativistic Shocks.
8. Xuening Bai's Lecture Notes Available on His Webpage.